

Single-phase circuits and electrical powers

1) Impedance and admittance

The impedance of a dipole D , through which a current $i(t)$ flows and which is subjected to a voltage $u(t)$ is define as the ratio:

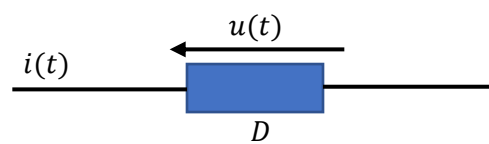
$$\underline{Z} = \frac{\underline{u}}{\underline{i}} = \frac{U}{I}$$


Figure 1: A dipole D

If:

$$u(t) = U_m \cos(\omega + \varphi) \text{ and } i(t) = I_m \cos(\omega t)$$

Then:

$$\underline{Z} = \frac{\underline{u}}{\underline{i}} = \frac{U}{I} = \frac{U_m e^{j\varphi}}{I_m e^{j0}} = Z e^{j\varphi} = Z \cos\varphi + jZ \sin\varphi = R + jX$$

Where :

- $R = Z \cos\varphi$: resistance in $[\Omega]$
- $X = Z \sin\varphi$: is the reactance in $[\Omega]$
- $|\underline{Z}| = Z = \frac{U_m}{I_m} = \sqrt{R^2 + X^2}$: is the module of $\underline{Z}$ in $[\Omega]$

The **admittance** of a dipole D is defined as the inverse of its impedance:

$$\underline{Y} = \frac{1}{\underline{Z}} = \frac{R}{R^2 + X^2} - j \frac{X}{R^2 + X^2} = G + jB$$

Where :

- $G = \frac{R}{R^2 + X^2}$: conductance in $[\Omega^{-1}]$
- $B = -\frac{X}{R^2 + X^2}$: susceptance in $[\Omega^{-1}]$
- $|\underline{Y}| = Y = \frac{1}{Z} = \frac{1}{\sqrt{R^2 + X^2}}$: magnitude of the admittance in $[\Omega^{-1}]$

2) Elements of an electrical circuit

A. Resistor R

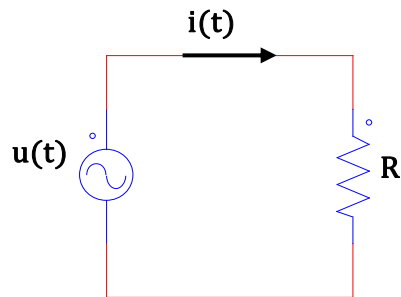


Figure 2: Circuit of a resistor R

Instantaneous relationship:

$$u(t) = Ri(t)$$

In the sinusoidal steady-state this relationship becomes in complex form:

$$\underline{u} = R\underline{i}$$

Hence:

$$\underline{Z}_R = \frac{\underline{u}}{\underline{i}} = R$$

With:

$$|\underline{Z}_R| = R \text{ and } \text{Arg}(\underline{Z}_R) = 0 \text{ rad}$$

If $i(t)$ is chosen as the phase reference (the phase origin):

- **Vector representation:** voltage u and current i are in phase.
- **Graphical representation:** voltage u and current i are in phase.

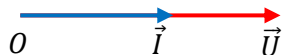


Figure 3: Vector representation

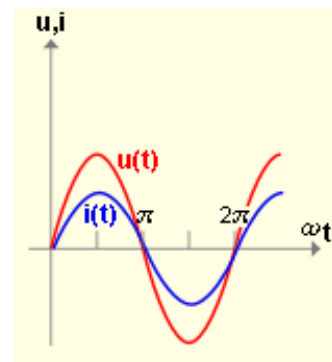


Figure 4: Graphical representation

B. Ideal Inductor L

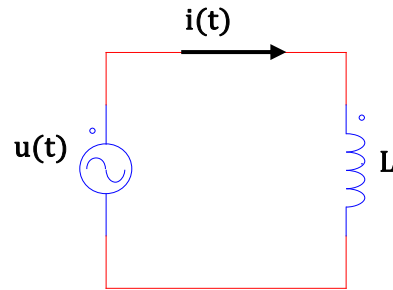


Figure 5: Circuit of an ideal inductor L

Instantaneous relationship:

$$u(t) = L \frac{di(t)}{dt}$$

In the sinusoidal steady-state this relationship becomes in complex form:

$$\underline{u} = jL\omega \underline{i}$$

Hence:

$$\underline{Z}_L = \frac{\underline{u}}{\underline{i}} = jL\omega$$

with:

$$|\underline{Z}_L| = L\omega \quad \text{And} \quad \text{Arg}(\underline{Z}_L) = \frac{\pi}{2} \text{ rad}$$

- **Vector representation:** the current i **lags** the voltage u by $\frac{\pi}{2}$ radians.
- **Graphical representation:** the voltage waveform u **leads** the current waveform i by a quarter period.

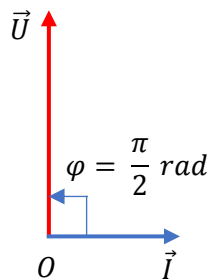


Figure 6: Vector representation

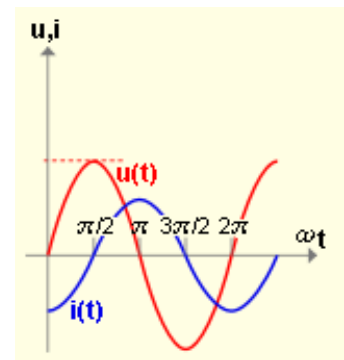


Figure 7: Graphical representation

C. Capacitor C

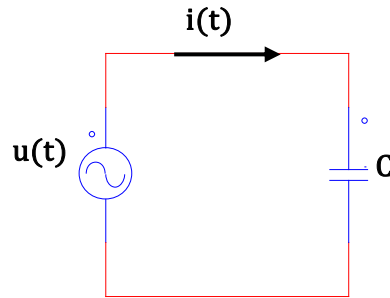


Figure 8: Circuit of a capacitor C

Instantaneous relationship:

$$u(t) = \frac{1}{C} \int i(t) dt$$

In the sinusoidal steady-state, the complex form is:

$$\underline{u} = \frac{1}{jC\omega} \underline{i}$$

Hence:

$$\underline{Z}_C = \frac{\underline{u}}{\underline{i}} = \frac{1}{jC\omega} = -j \frac{1}{C\omega}$$

With:

$$|\underline{Z}_C| = \frac{1}{C\omega} \text{ And } \text{Arg}(\underline{Z}_C) = -\frac{\pi}{2} \text{ rad}$$

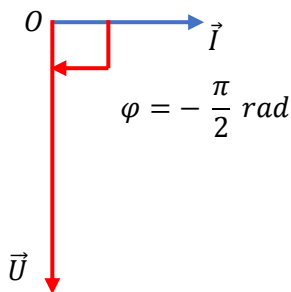


Figure 9: Vector representation

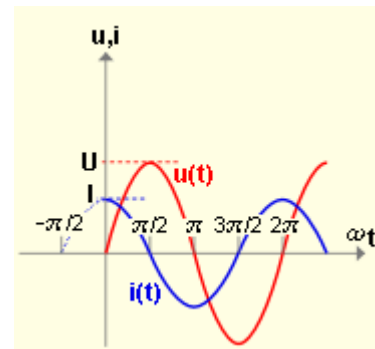


Figure 10: Graphical representation

- **Vector representation:** the current i leads the voltage u by $\frac{\pi}{2}$ radians.
- **Graphical representation:** the current waveform i reaches its maximum a quarter period before the voltage waveform u .

3) Fundamental laws of electricity

A. Generalized Ohm's law

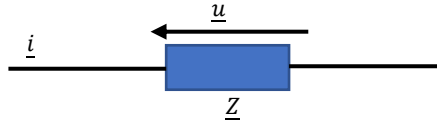


Figure11: A dipole D with $\underline{Z}$ impedance

The generalized Ohm's law in complex form relates the voltage $\underline{u}$ or $\underline{U}$ to the current $\underline{i}$ or $\underline{I}$ through the impedance $\underline{Z}$ which can be complex.

$$\underline{u} = \underline{z} \underline{i} \text{ ou } \underline{U} = \underline{Z} \underline{I}$$

This relationship holds in sinusoidal steady-state conditions where voltage and current are represented as complex numbers.

B. Kirchhoff's laws

- **Kirchhoff's voltage law (KVL)**

The phasor sum of the voltages around any closed loop is zero. This law arises from the conservation of energy in an electrical circuit.

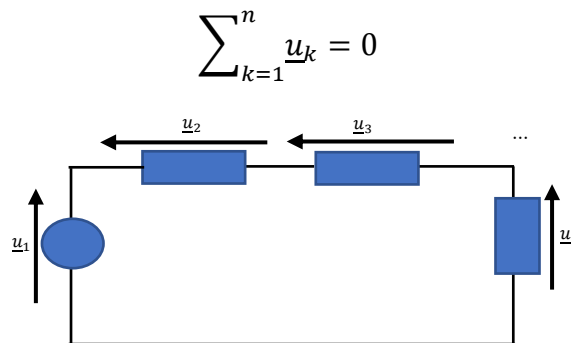


Figure 12: Electrical circuit with many dipoles

- **Kirchhoff's current law (KCL)**

The phasor sum of the currents entering and leaving any node is zero. This law arises from the conservation of electric charge.

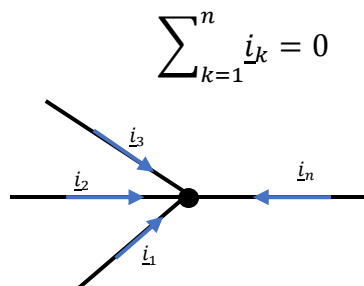


Figure 13: A node with many currents

- **Voltage and current divider**

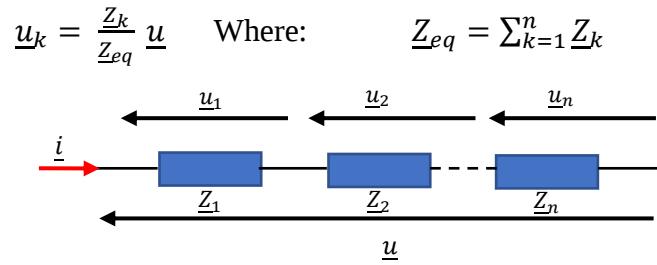


Figure 14: Voltage divider

And:
$$\underline{i}_k = \frac{\underline{Y}_k}{\sum_{k=1}^n \underline{Y}_k} \underline{i} \quad \text{Where:} \quad \sum_{k=1}^n \underline{Y}_k = \sum_{k=1}^n \frac{1}{\underline{Z}_k}$$

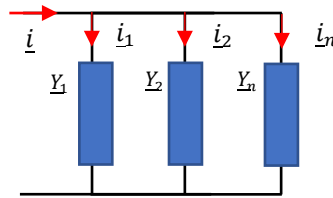


Figure15: Currents divider

4) Association of dipoles (impedances)

A. In series

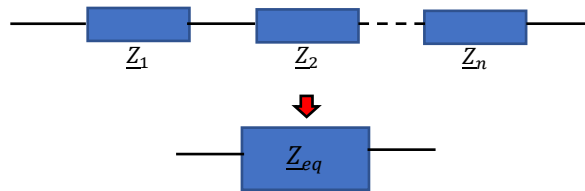


Figure16: Impedance equivalent in series

$$\underline{Z}_{eq} = \underline{Z}_1 + \underline{Z}_2 + \dots + \underline{Z}_n = \sum_{k=1}^n \underline{Z}_k$$

B. In parallel

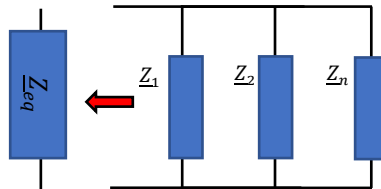


Figure17: Impedance equivalent in parallel

$$\frac{1}{\underline{Z}_{eq}} = \frac{1}{\underline{Z}_1} + \frac{1}{\underline{Z}_2} + \dots + \frac{1}{\underline{Z}_n} = \sum_{k=1}^n \frac{1}{\underline{Z}_k} = \sum_{k=1}^n \underline{Y}_k$$

C. Special cases

Elements	$\underline{Z}$	$ \underline{Z} $	$\varphi(\text{rad})$
R	R	R	0
L	$jL\omega$	$L\omega$	$\pi/2$
C	$-j/C\omega$	$1/C\omega$	$-\pi/2$
RL in series	$R + jL\omega$	$\sqrt{R^2 + L\omega^2}$	$\tan^{-1}(\frac{L\omega}{R})$
RC in series	$R - j\frac{1}{C\omega}$	$\sqrt{R^2 + (-\frac{1}{C\omega})^2}$	$\tan^{-1}(-\frac{1}{RC\omega})$
LC in series	$j(L\omega - \frac{1}{C\omega})$	$ L\omega - \frac{1}{C\omega} $	$\pm\pi/2$

5) Impedances triangle

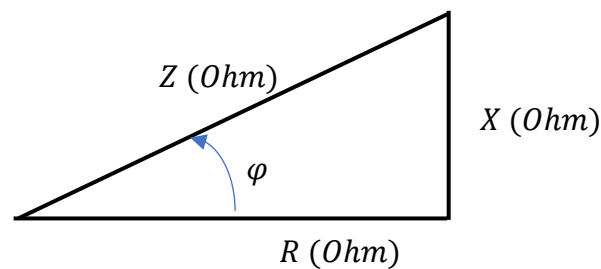


Figure 18: Impedances triangle

Electrical power in sinusoidal regime

1. Instantaneous power

We Consider a dipole D in receiver convention through which a current $i(t)$ flows and across which voltage $u(t)$ is applied.

The instantaneous power dissipated in the dipole D is expressed as:

$$P(t) = u(t) i(t)$$

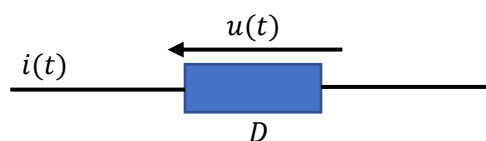


Figure 19: A dipole D

2. Instantaneous power in sinusoidal mode

For a sinusoidal steady-state the instantaneous power can be written as:

$$P(t) = U_{rms} I_{rms} \cos \varphi + U_{rms} I_{rms} \cos(2\omega t + \varphi)$$

- The first term represents the average (active) power.
- The second term the fluctuating power exchanged periodically between the source and the load.
- The fluctuating power (variable component of $P(t)$) is therefore:

$$P_f(t) = U_{rms} I_{rms} \cos(2\omega t + \varphi)$$

3. Active power

- The active power P is the average value of $P(t)$ over one period:

$$P = \frac{1}{T} \int_0^T P(t) dt = U_{rms} I_{rms} \cos \varphi \quad [\text{W}]$$

Note:

- For a pure resistor R

$$U_{rms} = R I_{rms} \Rightarrow P = R I_{rms}^2 \quad [\text{W}]$$

- For an ideal inductor L or capacitor C

$$\varphi = \pm \frac{\pi}{2} \text{ rad} \Rightarrow P = 0 \text{ rad} \quad [\text{W}]$$

4. Reactive power

The reactive power Q is the amplitude of the alternating component of $P(t)$:

$$Q = U_{rms} I_{rms} \sin \varphi \quad [\text{Var}]$$

Note:

- ❖ For a resistor R :

$$\varphi = 0 \Rightarrow Q = 0 \quad [\text{Var}]$$

- ❖ For an ideal inductor L :

$$U_{rms} = X_L I_{rms} \Rightarrow Q = X_L I_{rms}^2 \quad [\text{Var}]$$

Where: $X = L\omega$

- ❖ For a capacitor C :

$$Q = -X_C I_{rms}^2 \quad [\text{Var}]$$

Where :

$$X_C = \frac{1}{(C\omega)}$$

5. Apparent Power

The apparent power S represents the total power supplied by the source:

$$S = \sqrt{P^2 + Q^2} = U_{rms} I_{rms} = Z I_{rms}^2 \quad [\text{VA}]$$

It can be expressed in complex form:

$$\underline{S} = P + j Q$$

6. Power factor

The power factor P_F measures the efficiency of power transfer to the load

$$P_F = \cos \varphi = \frac{P}{S} \quad \text{Where : } 0 < P_F < 1$$

A low power factor indicates a large reactive component (inductive or capacitive load), leading to inefficient power transfer.

7. Power triangles

The power components can be represented geometrically in the power triangle

$$S^2 = P^2 + Q^2$$

P : Active power (horizontal axis)

Q : Reactive power (vertical axis)

S : Apparent power (hypotenuse)

φ : Phase shift between voltage and current

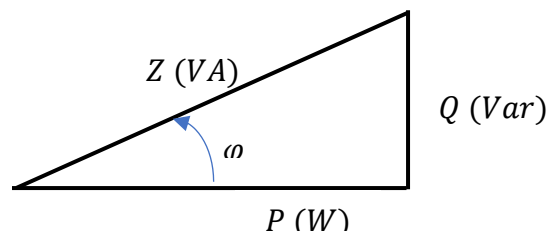


Figure 20: Powers triangle

8. Boucherot's Theorem

For a single-phase installation with several loads connected in parallel or series:

$$\diamond P_T = \sum_{k=1}^N P_k$$

$$\diamond Q_T = \sum_{k=1}^N Q_k$$

$$\diamond S_T \neq \sum_{k=1}^N S_k$$

However, in complex form:

$$\diamond \underline{S}_T = \sum_{k=1}^N \underline{S}_k$$

And the total apparent power is:

$$\diamond S_T = \sqrt{P_T^2 + Q_T^2}$$

The global power factor is then:

$$\diamond P_F = \cos \varphi_T = \frac{P_T}{S_T}$$

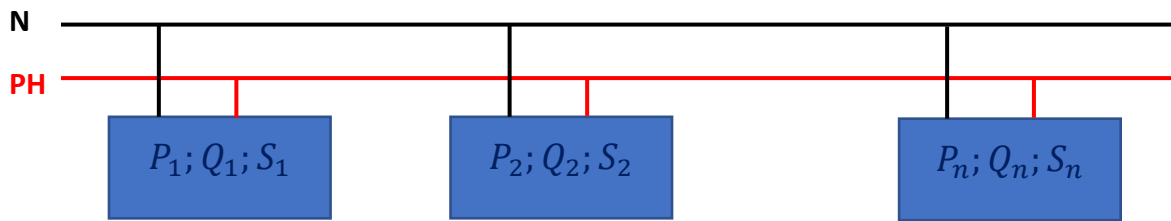


Figure 21 : Single phase installation